

# Spanning Trees on the Sierpinski Gasket

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*Received September 17, 2006; accepted November 17, 2006*  
*Published Online: January 2, 2007*

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We present the numbers of spanning trees on the Sierpinski gasket  $SG_d(n)$  at stage  $n$  with dimension  $d$  equal to two, three and four. The general expression for the number of spanning trees on  $SG_d(n)$  with arbitrary  $d$  is conjectured. The numbers of spanning trees on the generalized Sierpinski gasket  $SG_{d,b}(n)$  with  $d = 2$  and  $b = 3, 4$  are also obtained.

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**KEY WORDS:** Spanning trees, Sierpinski gasket, exact solutions

## 1. INTRODUCTION

The enumeration of the number of spanning trees  $N_{ST}(G)$  on a graph  $G$  was first considered by Kirchhoff in the analysis of electric circuits.<sup>(1)</sup> It is a problem of fundamental interest in mathematics<sup>(2–5)</sup> and physics.<sup>(6,7)</sup> The number of spanning trees corresponds to a special  $q \rightarrow 0$  limit of the partition function of the  $q$ -state Potts model in statistical mechanics,<sup>(8,9)</sup> which in turn is related to the sandpile model.<sup>(10)</sup> Just like other limits of the  $q$ -state Potts model, the spanning trees problem has been investigated intensely for decades, and has various applications in many areas. See, for example,<sup>(11)</sup> and references therein. It is also well known that there is a bijection between close-packed dimer coverings with spanning tree configurations on two related lattices.<sup>(12)</sup> Some recent studies on the enumeration of spanning trees and the calculation of their asymptotic growth constants on regular lattices were carried out in Refs. 13–16. It is of interest to consider spanning trees on self-similar fractal lattices which have scaling invariance

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rather than translational invariance. Fractals are geometric structures of noninteger Hausdorff dimension realized by repeated construction of an elementary shape on progressively smaller length scales.<sup>(17,18)</sup> A well-known example of fractal is the Sierpinski gasket which has been extensively studied in several contexts.<sup>(19–29)</sup> We shall derive rigorously the numbers of spanning trees on the Sierpinski gasket with dimension equal to two, three and four. The corresponding asymptotic growth constants have simple expressions. We shall also conjecture the general expression of the number of spanning trees on the Sierpinski gasket with arbitrary dimension.

## 2. PRELIMINARIES

We first recall some relevant definitions for spanning trees and the Sierpinski gasket in this section. A connected graph (without loops)  $G = (V, E)$  is defined by its vertex (site) and edge (bond) sets  $V$  and  $E$ .<sup>(2,30)</sup> Let  $v(G) = |V|$  be the number of vertices and  $e(G) = |E|$  the number of edges in  $G$ . A spanning subgraph  $G'$  is a subgraph of  $G$  with the same vertex set  $V$  and an edge set  $E' \subseteq E$ . As a tree is a connected graph with no circuits, a spanning tree on  $G$  is a spanning subgraph of  $G$  that is a tree and hence  $e(G') = v(G) - 1$ . The degree or coordination number  $k_i$  of a vertex  $v_i \in V$  is the number of edges attached to it. A  $k$ -regular graph is a graph with the property that each of its vertices has the same degree  $k$ . In general, one can associate an edge weight  $x_{ij}$  to each edge connecting adjacent vertices  $v_i$  and  $v_j$  (see, for example Ref. 13). For simplicity, all edge weights are set to one throughout this paper. A well-known method to enumerate spanning trees is to construct the Laplacian matrix  $Q(G)$ , which is the  $v(G) \times v(G)$  matrix with element  $Q(G)_{ij} = k_i \delta_{ij} - A(G)_{ij}$  where  $A(G)$  is the adjacency matrix. One of the eigenvalues of  $Q(G)$  is always zero. Denote the rest as  $\lambda(G)_i, 1 \leq i \leq v(G) - 1$ , then the number of spanning trees  $N_{ST}(G) = (1/v(G)) \prod_{i=1}^{v(G)-1} \lambda(G)_i$ .<sup>(2)</sup>

When the number of spanning trees  $N_{ST}(G)$  grows exponentially with  $v(G)$  as  $v(G) \rightarrow \infty$ , there exists a constant  $z_G$  describing this exponential growth:<sup>(4,5)</sup>

$$z_G = \lim_{v(G) \rightarrow \infty} \frac{\ln N_{ST}(G)}{v(G)} \quad (2.1)$$

where  $G$ , when used as a subscript in this manner, implicitly refers to the thermodynamic limit.

For a  $k$ -regular graph  $G_k$  with  $k \geq 3$ , there is an upper bound for the number of spanning trees<sup>(31,32)</sup>

$$N_{ST}(G_k) \leq \left( \frac{2 \ln v(G_k)}{v(G_k) k \ln k} \right) (b_k)^{v(G_k)}, \quad (2.2)$$

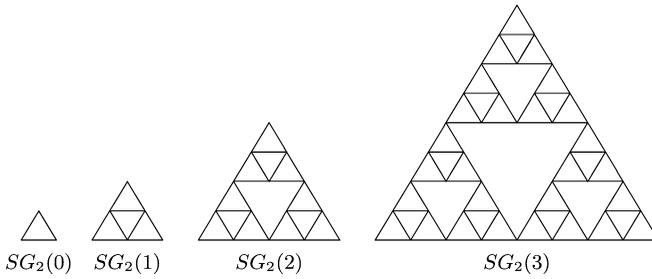


Fig. 1. The first four stages  $n = 0, 1, 2, 3$  of the two-dimensional Sierpinski gasket  $SG_2(n)$ .

where

$$b_k = \frac{(k - 1)^{k-1}}{[k(k - 2)]^{\frac{k}{2}-1}}. \tag{2.3}$$

With Eq. (2.1), this then yields<sup>(14)</sup>

$$z_{G_k} \leq \ln(b_k). \tag{2.4}$$

The construction of the two-dimensional Sierpinski gasket  $SG_2(n)$  at stage  $n$  is shown in Fig. 1. At stage  $n = 0$ , it is an equilateral triangle; while stage  $n + 1$  is obtained by the juxtaposition of three  $n$ -stage structures. In general, the Sierpinski gaskets  $SG_d$  can be built in any Euclidean dimension  $d$  with fractal dimensionality  $D = \ln(d + 1)/\ln 2$ .<sup>(20)</sup> For the Sierpinski gasket  $SG_d(n)$ , the numbers of edges and vertices are given by

$$e(SG_d(n)) = \binom{d + 1}{2} (d + 1)^n = \frac{d}{2} (d + 1)^{n+1}, \tag{2.5}$$

$$v(SG_d(n)) = \frac{d + 1}{2} [(d + 1)^n + 1]. \tag{2.6}$$

Except the  $(d + 1)$  outmost vertices which have degree  $d$ , all other vertices of  $SG_d(n)$  have degree  $2d$ . In the large  $n$  limit,  $SG_d$  is  $2d$ -regular.

The Sierpinski gasket can be generalized, denoted as  $SG_{d,b}(n)$ , by introducing the side length  $b$  which is an integer larger or equal to two.<sup>(33)</sup> The generalized Sierpinski gasket at stage  $n + 1$  is constructed with  $b$  layers of stage  $n$  hypertetrahedrons. The two-dimensional  $SG_{2,b}(n)$  with  $b = 3$  at stage  $n = 1, 2$  and  $b = 4$  at stage  $n = 1$  are illustrated in Fig. 2. The ordinary Sierpinski gasket  $SG_d(n)$  corresponds to the  $b = 2$  case, where the index  $b$  is neglected for simplicity. The Hausdorff dimension for  $SG_{d,b}$  is given by  $D = \ln\binom{b+d-1}{d} / \ln b$ .<sup>(33)</sup> Notice that  $SG_{d,b}$  is not  $k$ -regular even in the thermodynamic limit.

### 3. THE NUMBER OF SPANNING TREES ON $SG_2(n)$

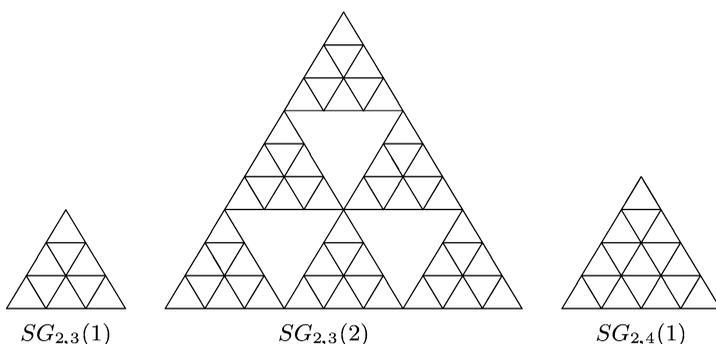
In this section we derive the number of spanning trees on the two-dimensional Sierpinski gasket  $SG_2(n)$  in detail. Let us start with the definitions of the quantities to be used.

**Definition 3.1.** Consider the two-dimensional Sierpinski gasket  $SG_2(n)$  at stage  $n$ . (a) Define  $f_2(n) \equiv N_{ST}(SG_2(n))$  as the number of spanning trees. (b) Define  $ga_2(n)$ ,  $gb_2(n)$ ,  $gc_2(n)$  as the number of spanning subgraphs with two trees such that one of the outmost vertices belongs to one tree and the other two outmost vertices belong to the other tree as illustrated in Fig. 3. (c) Define  $h_2(n)$  as the number of spanning subgraphs with three trees such that each of the outmost vertices belongs to a different tree.

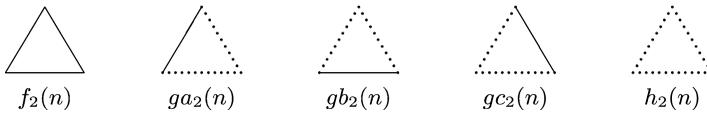
The quantities  $f_2(n)$ ,  $ga_2(n)$ ,  $gb_2(n)$ ,  $gc_2(n)$  and  $h_2(n)$  are illustrated in Fig. 3, where only the outmost vertices are shown. It is clear that the values  $ga_2(n)$ ,  $gb_2(n)$ ,  $gc_2(n)$  are the same because of rotation symmetry, and we define  $g_2(n) \equiv ga_2(n) = gb_2(n) = gc_2(n)$ . The initial values at stage 0 are  $f_2(0) = 3$ ,  $g_2(0) = 1$ ,  $h_2(0) = 1$ . Notice that  $h_2(n)$  is the number of spanning trees on  $SG_2(n)$  with the three outmost vertices identified. The purpose of this section is to obtain the expression of  $f_2(n)$  as follows.

**Theorem 3.1.** The number of spanning trees on the two-dimensional Sierpinski gasket  $SG_2(n)$  at stage  $n$  is given by

$$f_2(n) = 2^{\alpha_2(n)} 3^{\beta_2(n)} 5^{\gamma_2(n)}, \quad (3.1)$$



**Fig. 2.** The generalized two-dimensional Sierpinski gasket  $SG_{2,b}(n)$  with  $b = 3$  at stage  $n = 1, 2$  and  $b = 4$  at stage  $n = 1$ .



**Fig. 3.** Illustration for the spanning subgraphs  $f_2(n)$ ,  $ga_2(n)$ ,  $gb_2(n)$ ,  $gc_2(n)$  and  $h_2(n)$ . The two outmost vertices at the ends of a solid line belong to one tree, while the two outmost vertices at the ends of a dot line belong to separated trees.

where the exponents are

$$\alpha_2(n) = \frac{1}{2}(3^n - 1), \tag{3.2}$$

$$\beta_2(n) = \frac{1}{4}(3^{n+1} + 2n + 1), \tag{3.3}$$

$$\gamma_2(n) = \frac{1}{4}(3^n - 2n - 1). \tag{3.4}$$

This theorem can be proved by the following two lemmas. The three quantities  $f_2(n)$ ,  $g_2(n)$  and  $h_2(n)$  satisfy recursion relations which were first obtained in Ref. 27.

**Lemma 1.** For any non-negative integer  $n$ ,

$$f_2(n + 1) = 6f_2^2(n)g_2(n), \tag{3.5}$$

$$g_2(n + 1) = f_2^2(n)h_2(n) + 7f_2(n)g_2^2(n), \tag{3.6}$$

$$h_2(n + 1) = 12f_2(n)g_2(n)h_2(n) + 14g_2^3(n). \tag{3.7}$$

**Proof:** The Sierpinski gaskets  $SG_2(n + 1)$  is composed of three  $SG_2(n)$  with three pairs of vertices identified. To obtain the number of spanning trees  $f_2(n + 1)$ , one of the  $SG_2(n)$  should be spanned by two trees. There are six possibilities as illustrated in Fig. 4. Therefore, we have

$$f_2(n + 1) = 2f_2^2(n)[ga_2(n) + gb_2(n) + gc_2(n)] = 6f_2^2(n)g_2(n). \tag{3.8}$$

Similarly,  $ga_2(n + 1)$  for  $SG_2(n + 1)$  can be obtained with appropriated configurations of its three constituting  $SG_2(n)$  as illustrated in Fig. 5. Thus,

$$ga_2(n + 1) = f_2^2(n)h_2(n) + 3f_2(n)ga_2^2(n) + 2f_2(n)ga_2(n)gc_2(n) + 2f_2(n)ga_2(n)gb_2(n). \tag{3.9}$$

With the identity  $ga_2(n) = gb_2(n) = gc_2(n) = g_2(n)$ , Eq. (3.6) is verified.



Fig. 4. Illustration for the expression of  $f_2(n + 1)$ .

Finally,  $h_2(n + 1)$  is the summation of appropriated configurations as illustrated in Fig. 6, so that

$$\begin{aligned}
 h_2(n + 1) &= 4f_2(n)h_2(n)[ga_2(n) + gb_2(n) + gc_2(n)] \\
 &\quad + 2gc_2(n)ga_2(n)[gc_2(n) + ga_2(n)] \\
 &\quad + 2ga_2(n)gb_2(n)[ga_2(n) + gb_2(n)] \\
 &\quad + 2gb_2(n)gc_2(n)[gb_2(n) + gc_2(n)] \\
 &\quad + 2ga_2(n)gb_2(n)gc_2(n). \tag{3.10}
 \end{aligned}$$

With the identity  $ga_2(n) = gb_2(n) = gc_2(n) = g_2(n)$ , Eq. (3.7) is verified.  $\square$

By the three Eqs. (3.5)–(3.7),  $f_2(n)$ ,  $g_2(n)$  and  $h_2(n)$  can be solved.

**Lemma 2.** For any non-negative integer  $n$ ,

$$f_2(n) = 2^{\alpha_2(n)} 3^{\beta_2(n)} 5^{\gamma_2(n)}, \tag{3.11}$$

$$g_2(n) = 2^{\alpha_2(n)} 3^{\beta_2(n)-n-1} 5^{\gamma_2(n)+n}, \tag{3.12}$$

$$h_2(n) = 2^{\alpha_2(n)} 3^{\beta_2(n)-2n-1} 5^{\gamma_2(n)+2n} \tag{3.13}$$

where  $\alpha_2(n)$ ,  $\beta_2(n)$  and  $\gamma_2(n)$  are given in Eqs. (3.2)–(3.4)

**Proof:** As  $\alpha_2(0) = 0$ ,  $\beta_2(0) = 1$  and  $\gamma_2(0) = 0$ , Eqs. (3.11)–(3.13) are correct for  $n = 0$ . By the recursion relation (3.5), it is easy to obtain

$$f_2(n + 1) = 2^{3\alpha_2(n)+1} 3^{3\beta_2(n)-n} 5^{3\gamma_2(n)+n} \tag{3.14}$$

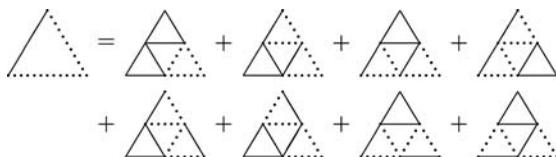


Fig. 5. Illustration for the expression of  $ga_2(n + 1)$ .

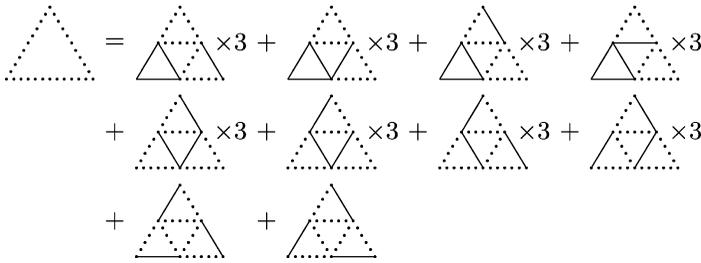


Fig. 6. Illustration for the expression of  $h_2(n + 1)$ . The multiplication for the eight configurations on the right-hand-side corresponds to three possible orientations.

where

$$3\alpha_2(n) + 1 = \frac{1}{2}(3^{n+1} - 1) = \alpha_2(n + 1), \tag{3.15}$$

$$3\beta_2(n) - n = \frac{1}{4}(3^{n+2} + 2n + 3) = \beta_2(n + 1), \tag{3.16}$$

$$3\gamma_2(n) + n = \frac{1}{4}(3^{n+1} - 2n - 3) = \gamma_2(n + 1). \tag{3.17}$$

The proof for Eq. (3.11) is completed by induction. Similarly, we have

$$g_2(n + 1) = 2^{3\alpha_2(n)+1} 3^{3\beta_2(n)-2n-2} 5^{3\gamma_2(n)+2n+1}, \tag{3.18}$$

$$h_2(n + 1) = 2^{3\alpha_2(n)+1} 3^{3\beta_2(n)-3n-3} 5^{3\gamma_2(n)+3n+2}, \tag{3.19}$$

which verify Eqs. (3.12) and (3.13). □

We note that although the Laplacian matrix  $Q(G)$  for  $G = SG_2(n)$  can be constructed, it does not look simple to diagonalize. For example, with appropriate order of the vertices,

$$Q(SG_2(1)) = \begin{pmatrix} 4 & -1 & -1 & 0 & -1 & -1 \\ -1 & 4 & -1 & -1 & 0 & -1 \\ -1 & -1 & 4 & -1 & -1 & 0 \\ 0 & -1 & -1 & 2 & 0 & 0 \\ -1 & 0 & -1 & 0 & 2 & 0 \\ -1 & -1 & 0 & 0 & 0 & 2 \end{pmatrix} \tag{3.20}$$

and

$Q(SG_2(2))$

$$= \begin{pmatrix} 4 & -1 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & -1 & 0 & 0 & 0 \\ -1 & 4 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & -1 & 0 & 0 \\ -1 & -1 & 4 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & -1 & 0 & 0 \\ 0 & 0 & 0 & 4 & -1 & -1 & 0 & 0 & 0 & -1 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 4 & -1 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & -1 \\ 0 & 0 & 0 & -1 & -1 & 4 & 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 4 & -1 & -1 & -1 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 4 & -1 & -1 & 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & -1 & 4 & 0 & -1 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 & 0 & -1 & -1 & -1 & 0 & 4 & 0 & 0 & 0 & 0 & 0 \\ -1 & -1 & 0 & 0 & 0 & 0 & -1 & 0 & -1 & 0 & 4 & 0 & 0 & 0 & 0 \\ -1 & 0 & -1 & -1 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 4 & 0 & 0 & 0 \\ 0 & -1 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & -1 & 0 & 0 & 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & 0 & -1 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 \end{pmatrix}. \tag{3.21}$$

Denote the non-zero eigenvalues of  $Q(SG_2(n))$  as  $\lambda(SG_2(n))_i$  for  $1 \leq i \leq v(SG_2(n)) - 1$ .

**Corollary 3.1.** *The product of non-zero eigenvalues of  $Q(SG_2(n))$  is given by*

$$\prod_{i=1}^{v(SG_2(n))-1} \lambda(SG_2(n))_i = v(SG_2(n))f_2(n), \tag{3.22}$$

where  $f_2(n)$  is given in Theorem 3.1 and the number of vertices  $v(SG_2(n)) = 3/2(3^n + 1)$  for  $SG_2(n)$  is given by Eq. (2.6) with  $d = 2$ . Similar expressions apply to all other cases discussed below.

By the definition in Eq. (2.1), we have the following corollary.

**Corollary 3.2.** *The asymptotic growth constant for  $SG_2$  is given by*

$$\begin{aligned} z_{SG_2} &= \frac{1}{3} \ln 2 + \frac{1}{2} \ln 3 + \frac{1}{6} \ln 5 \\ &\simeq 1.048594856 \dots \end{aligned} \tag{3.23}$$

This is equivalent to Eq. (20) of Ref. 28. In passing, we notice that the number of spanning trees is the same for a planar graph and its dual, so that Theorem 3.1 also applies to the dual of  $SG_2(n)$ , denoted as  $SG_2^*(n)$ . As  $SG_2$  is 4-regular in the large  $n$  limit, we have  $z_{SG_2^*} = z_{SG_2}$ .<sup>(16)</sup>

#### 4. THE NUMBER OF SPANNING TREES ON $SG_{2,b}(n)$ WITH $b = 3, 4$

The method given in the previous section can be applied to the number of spanning trees on  $SG_{d,b}(n)$  with larger values of  $d$  and  $b$ . The number of configurations to be considered increases as  $d$  and  $b$  increase, and the recursion relations must be derived individually for each  $d$  and  $b$ . In this section, we consider the generalized two-dimensional Sierpinski gasket  $SG_{2,b}(n)$  with the number of layers  $b$  equal to three and four. For  $SG_{2,3}(n)$ , the numbers of edges and vertices are given by

$$e(SG_{2,3}(n)) = 3 \times 6^n, \tag{4.1}$$

$$v(SG_{2,3}(n)) = \frac{7 \times 6^n + 8}{5}, \tag{4.2}$$

where the three outmost vertices have degree two. There are  $(6^n - 1)/5$  vertices of  $SG_{2,3}(n)$  with degree six and  $6(6^n - 1)/5$  vertices with degree four. Define  $f_{2,3}(n)$ ,  $g_{2,3}(n)$ ,  $h_{2,3}(n)$  as in Definition 3.1 such that the initial values are  $f_{2,3}(0) = 3$ ,  $g_{2,3}(0) = 1$ ,  $h_{2,3}(0) = 1$ . By the method illustrated in the above section, we obtain following recursion relations for any non-negative integer  $n$ .

$$f_{2,3}(n + 1) = 18f_{2,3}^4(n)g_{2,3}(n)h_{2,3}(n) + 142f_{2,3}^3(n)g_{2,3}^3(n), \tag{4.3}$$

$$g_{2,3}(n + 1) = 2f_{2,3}^4(n)h_{2,3}^2(n) + 77f_{2,3}^3(n)g_{2,3}^2(n)h_{2,3}(n) + 171f_{2,3}^2(n)g_{2,3}^4(n), \tag{4.4}$$

$$h_{2,3}(n + 1) = 60f_{2,3}^3(n)g_{2,3}(n)h_{2,3}^2(n) + 564f_{2,3}^2(n)g_{2,3}^3(n)h_{2,3}(n) + 468f_{2,3}(n)g_{2,3}^5(n). \tag{4.5}$$

The figures for these configurations are too many to be shown here. By induction as in Lemma 3.3, these equations can be solved.

**Theorem 4.1.** *The number of spanning trees on the two-dimensional Sierpinski gasket  $SG_{2,3}(n)$  at stage  $n$  is given by*

$$f_{2,3}(n) = 2^{\alpha_{2,3}(n)} 3^{\beta_{2,3}(n)} 5^{\gamma_{2,3}(n)} 7^{\delta_{2,3}(n)}, \tag{4.6}$$

where the exponents are

$$\alpha_{2,3}(n) = \frac{2}{5}(6^n - 1), \tag{4.7}$$

$$\beta_{2,3}(n) = \frac{1}{25}(13 \times 6^n - 15n + 12), \quad (4.8)$$

$$\gamma_{2,3}(n) = \frac{1}{25}(3 \times 6^n - 15n - 3), \quad (4.9)$$

$$\delta_{2,3}(n) = \frac{1}{25}(7 \times 6^n + 15n - 7). \quad (4.10)$$

$g_{2,3}(n)$  and  $h_{2,3}(n)$  can also be expressed by these exponents:

$$\begin{aligned} g_{2,3}(n) &= 2^{\alpha_{2,3}(n)} 3^{\beta_{2,3}(n)+n-1} 5^{\gamma_{2,3}(n)+n} 7^{\delta_{2,3}(n)-n}, \\ h_{2,3}(n) &= 2^{\alpha_{2,3}(n)} 3^{\beta_{2,3}(n)+2n-1} 5^{\gamma_{2,3}(n)+2n} 7^{\delta_{2,3}(n)-2n}. \end{aligned} \quad (4.11)$$

By the definition in Eq. (2.1), we have the following corollary.

**Corollary 4.1.** *The asymptotic growth constant for  $SG_{2,3}$  is given by*

$$\begin{aligned} z_{SG_{2,3}} &= \frac{2}{7} \ln 2 + \frac{13}{35} \ln 3 + \frac{3}{35} \ln 5 + \frac{1}{5} \ln 7 \\ &\simeq 1.133231895 \dots \end{aligned} \quad (4.12)$$

For  $SG_{2,4}(n)$ , the numbers of edges and vertices are given by

$$e(SG_{2,4}(n)) = 3 \times 10^n, \quad (4.13)$$

$$v(SG_{2,4}(n)) = \frac{4 \times 10^n + 5}{3}, \quad (4.14)$$

where again the three outmost vertices have degree two. There are  $(10^n - 1)/3$  vertices of  $SG_{2,4}(n)$  with degree six, and  $(10^n - 1)$  vertices with degree four. Define  $f_{2,4}(n)$ ,  $g_{2,4}(n)$ ,  $h_{2,4}(n)$  as in Definition 3.1 such that the initial values are  $f_{2,4}(0) = 3$ ,  $g_{2,4}(0) = 1$ ,  $h_{2,4}(0) = 1$ . We write a computer program to obtain following recursion relations for any non-negative integer  $n$ .

$$\begin{aligned} f_{2,4}(n+1) &= 2f_{2,4}^7(n)h_{2,4}^3(n) + 516f_{2,4}^6(n)g_{2,4}^2(n)h_{2,4}^2(n) \\ &\quad + 5856f_{2,4}^5(n)g_{2,4}^4(n)h_{2,4}(n) + 11354f_{2,4}^4(n)g_{2,4}^6(n), \end{aligned} \quad (4.15)$$

$$\begin{aligned} g_{2,4}(n+1) &= 82f_{2,4}^6(n)g_{2,4}(n)h_{2,4}^3(n) + 2786f_{2,4}^5(n)g_{2,4}^3(n)h_{2,4}^2(n) \\ &\quad + 14480f_{2,4}^4(n)g_{2,4}^5(n)h_{2,4}(n) + 13732f_{2,4}^3(n)g_{2,4}^7(n), \end{aligned} \quad (4.16)$$

$$\begin{aligned}
 h_{2,4}(n+1) &= 20f_{2,4}^6(n)h_{2,4}^4(n) + 2388f_{2,4}^5(n)g_{2,4}^2(n)h_{2,4}^3(n) \\
 &\quad + 30948f_{2,4}^4(n)g_{2,4}^4(n)h_{2,4}^2(n) + 83234f_{2,4}^3(n)g_{2,4}^6(n)h_{2,4}(n) \\
 &\quad + 42210f_{2,4}^2(n)g_{2,4}^8(n).
 \end{aligned}
 \tag{4.17}$$

By induction as in Lemma 3.3, these equations can be solved.

**Theorem 4.2.** *The number of spanning trees on the two-dimensional Sierpinski gasket  $SG_{2,4}(n)$  at stage  $n$  is given by*

$$f_{2,4}(n) = 2^{\alpha_{2,4}(n)}3^{\beta_{2,4}(n)}5^{\gamma_{2,4}(n)}41^{\delta_{2,4}(n)}103^{\epsilon_{2,4}(n)},
 \tag{4.18}$$

where the exponents are

$$\alpha_{2,4}(n) = \frac{2}{9}(10^n - 1),
 \tag{4.19}$$

$$\beta_{2,4}(n) = \frac{1}{3}(10^n + 2),
 \tag{4.20}$$

$$\gamma_{2,4}(n) = \frac{1}{9}(10^n - 1),
 \tag{4.21}$$

$$\delta_{2,4}(n) = \frac{2}{27}(2 \times 10^n + 9n - 2),
 \tag{4.22}$$

$$\epsilon_{2,4}(n) = \frac{2}{27}(10^n - 9n - 1).
 \tag{4.23}$$

$g_{2,4}(n)$  and  $h_{2,4}(n)$  can also be expressed by these exponents:

$$\begin{aligned}
 g_{2,4}(n) &= 2^{\alpha_{2,4}(n)}3^{\beta_{2,4}(n)-1}5^{\gamma_{2,4}(n)}41^{\delta_{2,4}(n)-n}103^{\epsilon_{2,4}(n)+n}, \\
 h_{2,4}(n) &= 2^{\alpha_{2,4}(n)}3^{\beta_{2,4}(n)-1}5^{\gamma_{2,4}(n)}41^{\delta_{2,4}(n)-2n}103^{\epsilon_{2,4}(n)+2n}.
 \end{aligned}
 \tag{4.24}$$

By the definition in Eq. (2.1), we have the following corollary.

**Corollary 4.2.** *The asymptotic growth constant for  $SG_{2,4}$  is given by*

$$\begin{aligned}
 z_{SG_{2,4}} &= \frac{1}{6} \ln 2 + \frac{1}{4} \ln 3 + \frac{1}{12} \ln 5 + \frac{1}{9} \ln(41) + \frac{1}{18} \ln(103) \\
 &\simeq 1.194401490\dots
 \end{aligned}
 \tag{4.25}$$

### 5. THE NUMBER OF SPANNING TREES ON $SG_d(n)$ WITH $d = 3, 4$

In this section, we derive the number of spanning trees on  $SG_d(n)$  with  $d = 3, 4$ . For the three-dimensional Sierpinski gasket  $SG_3(n)$ , we use the following definitions.

**Definition 5.1.** Consider the three-dimensional Sierpinski gasket  $SG_3(n)$  at stage  $n$ . (a) Define  $f_3(n) \equiv N_{ST}(SG_3(n))$  as the number of spanning trees. (b) Define  $g_3(n)$  as the number of spanning subgraphs with two trees such that one of the outmost vertices belongs to one tree and the other three outmost vertices belong to the other tree. (c) Define  $h_3(n)$  as the number of spanning subgraphs with two trees such that two of the outmost vertices belong to one tree and the other two outmost vertices belong to the other tree. (d) Define  $p_3(n)$  as the number of spanning subgraphs with three trees such that two of the outmost vertices belong to one tree and the other two outmost vertices separately belong to the other trees. (e) Define  $q_3(n)$  as the number of spanning subgraphs with four trees such that each of the outmost vertices belongs to a different tree.

The quantities  $f_3(n)$ ,  $g_3(n)$ ,  $h_3(n)$ ,  $p_3(n)$  and  $q_3(n)$  are illustrated in Fig. 7, where only the outmost vertices are shown. There are four equivalent  $g_3(n)$ , three equivalent  $h_3(n)$ , and six equivalent  $p_3(n)$ . The initial values at stage 0 are  $f_3(0) = 16$ ,  $g_3(0) = 3$ ,  $h_3(0) = 1$ ,  $p_3(0) = 1$ ,  $q_3(0) = 1$ . Notice that  $q_3(n)$  is the number of spanning trees on  $SG_3(n)$  with the four outmost vertices identified.

In the process of calculation, we find that it is convenient to combine  $g_3(n)$  and  $h_3(n)$  and define  $gh_3(n) \equiv g_3(n) + h_3(n)$ . We obtain following recursion relations for any non-negative integer  $n$ .

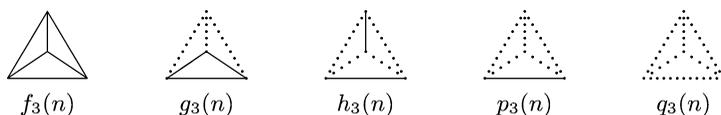
$$f_3(n+1) = 72f_3^2(n)gh_3(n)p_3(n) + 56f_3(n)gh_3^3(n), \quad (5.1)$$

$$gh_3(n+1) = 6f_3^2(n)gh_3(n)q_3(n) + 26f_3^2(n)p_3^2(n) + 120f_3(n)gh_3^2(n)p_3(n) + 22gh_3^4(n), \quad (5.2)$$

$$p_3(n+1) = 6f_3^2(n)p_3(n)q_3(n) + 14f_3(n)gh_3^2(n)q_3(n) + 120f_3(n)gh_3(n)p_3^2(n) + 88gh_3^3(n)p_3(n), \quad (5.3)$$

$$q_3(n+1) = 144f_3(n)gh_3(n)p_3(n)q_3(n) + 208f_3(n)p_3^3(n) + 56gh_3^3(n)q_3(n) + 720gh_3^2(n)p_3^2(n). \quad (5.4)$$

The figures for these configurations are too many to be shown here. With the initial values given above, these equations can be solved by induction as in Lemma 3.3.



**Fig. 7.** Illustration for the spanning subgraphs  $f_3(n)$ ,  $g_3(n)$ ,  $h_3(n)$ ,  $p_3(n)$  and  $q_3(n)$ . The two outmost vertices at the ends of a solid line belong to one tree, while the two outmost vertices at the ends of a dot line belong to separated trees.

**Theorem 5.1.** *The number of spanning trees on the three-dimensional Sierpinski gasket  $SG_3(n)$  at stage  $n$  is given by*

$$f_3(n) = 2^{\alpha_3(n)} 3^{\beta_3(n)}, \quad (5.5)$$

where the exponents are

$$\alpha_3(n) = 4^{n+1} + n, \quad (5.6)$$

$$\beta_3(n) = \frac{1}{3}(4^n - 3n - 1). \quad (5.7)$$

$gh_3(n)$ ,  $p_3(n)$  and  $q_3(n)$  can also be expressed by these exponents:

$$\begin{aligned} gh_3(n) &= 2^{\alpha_3(n)-n-2} 3^{\beta_3(n)+n}, \\ p_3(n) &= 2^{\alpha_3(n)-2n-4} 3^{\beta_3(n)+2n}, \\ q_3(n) &= 2^{\alpha_3(n)-3n-4} 3^{\beta_3(n)+3n}. \end{aligned} \quad (5.8)$$

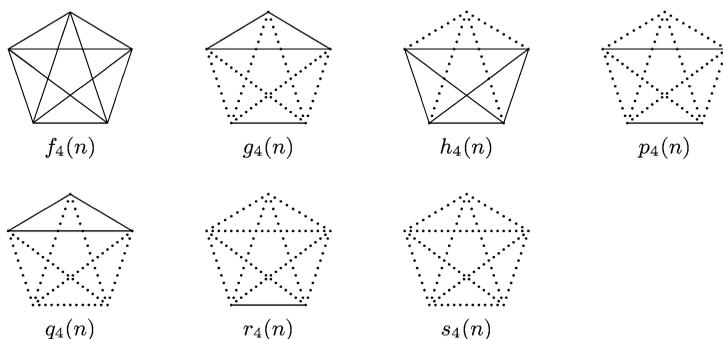
By the definition in Eq. (2.1), we have the following corollary.

**Corollary 5.1.** *The asymptotic growth constant for  $SG_3$  is given by*

$$\begin{aligned} z_{SG_3} &= 2 \ln 2 + \frac{1}{6} \ln 3 \\ &\simeq 1.569396409 \dots \end{aligned} \quad (5.9)$$

For the four-dimensional Sierpinski gasket  $SG_4(n)$ , we use the following definitions.

**Definition 5.2.** *Consider the four-dimensional Sierpinski gasket  $SG_4(n)$  at stage  $n$ . (a) Define  $f_4(n) \equiv N_{ST}(SG_4(n))$  as the number of spanning trees. (b) Define  $g_4(n)$  as the number of spanning subgraphs with two trees such that two of the outmost vertices belong to one tree and the other three outmost vertices belong to the other tree. (c) Define  $h_4(n)$  as the number of spanning subgraphs with two trees such that one of the outmost vertices belong to one tree and the other four outmost vertices belong to the other tree. (d) Define  $p_4(n)$  as the number of spanning subgraphs with three trees such that one of the outmost vertices belong to one tree, two of the other outmost vertices belong to another tree and the rest two outmost vertices belong to the other tree. (e) Define  $q_4(n)$  as the number of spanning subgraphs with three trees such that three of the outmost vertices belong to one tree and the other two outmost vertices separately belong to the other trees. (f) Define  $r_4(n)$  as the number of spanning subgraphs with four trees such*



**Fig. 8.** Illustration for the spanning subgraphs  $f_4(n)$ ,  $g_4(n)$ ,  $h_4(n)$ ,  $p_4(n)$ ,  $q_4(n)$ ,  $r_4(n)$  and  $s_4(n)$ . The two outmost vertices at the ends of a solid line belong to one tree, while the two outmost vertices at the ends of a dot line belong to separated trees.

that two of the outmost vertices belong to one tree and the other three outmost vertices separately belong to the other trees. (g) Define  $s_4(n)$  as the number of spanning subgraphs with five trees such that each of the outmost vertices belongs to a different tree.

The quantities  $f_4(n)$ ,  $g_4(n)$ ,  $h_4(n)$ ,  $p_4(n)$ ,  $q_4(n)$ ,  $r_4(n)$  and  $s_4(n)$  are illustrated in Fig. 8, where only the outmost vertices are shown. There are ten equivalent  $g_4(n)$ , five equivalent  $h_4(n)$ , fifteen equivalent  $p_4(n)$ , ten equivalent  $q_4(n)$  and ten equivalent  $r_4(n)$ . The initial values at stage 0 are  $f_4(0) = 125$ ,  $g_4(0) = 3$ ,  $h_4(0) = 16$ ,  $p_4(0) = 1$ ,  $q_4(0) = 3$ ,  $r_4(0) = 1$ ,  $s_4(0) = 1$ . Notice that  $s_4(n)$  is the number of spanning trees on  $SG_4(n)$  with the five outmost vertices identified.

We find that it is convenient to reduce the number of variables by defining  $gh_4(n) \equiv 3g_4(n) + h_4(n)$  and  $pq_4(n) \equiv 2p_4(n) + q_4(n)$ . We write a computer program to obtain following recursion relations for any non-negative integer  $n$ .

$$\begin{aligned}
 f_4(n+1) &= 1440f_4^2(n)gh_4(n)pq_4(n)r_4(n) + 520f_4^2(n)pq_4^3(n) \\
 &\quad + 1120f_4(n)gh_4^3(n)r_4(n) + 3600f_4(n)gh_4^2(n)pq_4^2(n) \\
 &\quad + 1320gh_4^4(n)pq_4(n), \tag{5.10}
 \end{aligned}$$

$$\begin{aligned}
 gh_4(n+1) &= 72f_4^2(n)gh_4(n)pq_4(n)s_4(n) + 378f_4^2(n)gh_4(n)r_4^2(n) \\
 &\quad + 816f_4^2(n)pq_4^2(n)r_4(n) + 56f_4(n)gh_4^3(n)s_4(n) \\
 &\quad + 3756f_4(n)gh_4^2(n)pq_4(n)r_4(n) + 2360f_4(n)gh_4(n)pq_4^3(n) \\
 &\quad + 688gh_4^4(n)r_4(n) + 2562gh_4^3(n)pq_4^2(n), \tag{5.11}
 \end{aligned}$$

$$\begin{aligned}
pq_4(n+1) &= 48f_4^2(n)gh_4(n)r_4(n)s_4(n) + 52f_4^2(n)pq_4^2(n)s_4(n) \\
&\quad + 544f_4^2(n)pq_4(n)r_4^2(n) + 240f_4(n)gh_4^2(n)pq_4(n)s_4(n) \\
&\quad + 1252f_4(n)gh_4^2(n)r_4^2(n) + 4720f_4(n)gh_4(n)pq_4^2(n)r_4(n) \\
&\quad + 724f_4(n)pq_4^4(n) + 44gh_4^4(n)s_4(n) \\
&\quad + 3416gh_4^3(n)pq_4(n)r_4(n) + 3104gh_4^2(n)pq_4^3(n), \quad (5.12)
\end{aligned}$$

$$\begin{aligned}
r_4(n+1) &= 72f_4^2(n)pq_4(n)r_4(n)s_4(n) + 126f_4^2(n)r_4^3(n) \\
&\quad + 168f_4(n)gh_4^2(n)r_4(n)s_4(n) + 360f_4(n)gh_4(n)pq_4^2(n)s_4(n) \\
&\quad + 3756f_4(n)gh_4(n)pq_4(n)r_4^2(n) + 2360f_4(n)pq_4^3(n)r_4(n) \\
&\quad + 264gh_4^3(n)pq_4(n)s_4(n) + 1376gh_4^3(n)r_4^2(n) \\
&\quad + 7686gh_4^2(n)pq_4^2(n)r_4(n) + 2328gh_4(n)pq_4^4(n), \quad (5.13)
\end{aligned}$$

$$\begin{aligned}
s_4(n+1) &= 2880f_4(n)gh_4(n)pq_4(n)r_4(n)s_4(n) + 5040f_4(n)gh_4(n)r_4^3(n) \\
&\quad + 1040f_4(n)pq_4^3(n)s_4(n) + 16320f_4(n)pq_4^2(n)r_4^2(n) \\
&\quad + 1120gh_4^3(n)r_4(n)s_4(n) + 3600gh_4^2(n)pq_4^2(n)s_4(n) \\
&\quad + 37560gh_4^2(n)pq_4(n)r_4^2(n) + 47200gh_4(n)pq_4^3(n)r_4(n) \\
&\quad + 4344pq_4^5(n). \quad (5.14)
\end{aligned}$$

With the initial values given above, these equations can be solved by induction as in Lemma 3.3.

**Theorem 5.2.** *The number of spanning trees on the four-dimensional Sierpinski gasket  $SG_4(n)$  at stage  $n$  is given by*

$$f_4(n) = 2^{\alpha_4(n)} 5^{\beta_4(n)} 7^{\gamma_4(n)}, \quad (5.15)$$

where the exponents are

$$\alpha_4(n) = \frac{3}{2}(5^n - 1), \quad (5.16)$$

$$\beta_4(n) = \frac{3}{8}(5^{n+1} + 4n + 3), \quad (5.17)$$

$$\gamma_4(n) = \frac{3}{8}(5^n - 4n - 1). \quad (5.18)$$

$gh_4(n)$ ,  $pq_4(n)$ ,  $r_4(n)$  and  $s_4(n)$  can also be expressed by these exponents:

$$\begin{aligned} gh_4(n) &= 2^{\alpha_4(n)} 5^{\beta_4(n)-n-1} 7^{\gamma_4(n)+n}, \\ pq_4(n) &= 2^{\alpha_4(n)} 5^{\beta_4(n)-2n-2} 7^{\gamma_4(n)+2n}, \\ r_4(n) &= 2^{\alpha_4(n)} 5^{\beta_4(n)-3n-3} 7^{\gamma_4(n)+3n}, \\ s_4(n) &= 2^{\alpha_4(n)} 5^{\beta_4(n)-4n-3} 7^{\gamma_4(n)+4n}. \end{aligned} \quad (5.19)$$

By the definition in Eq. (2.1), we have the following corollary.

**Corollary 5.2.** *The asymptotic growth constant for  $SG_4$  is given by*

$$\begin{aligned} z_{SG_4} &= \frac{3}{5} \ln 2 + \frac{3}{4} \ln 5 + \frac{3}{20} \ln 7 \\ &\simeq 1.914853265 \dots \end{aligned} \quad (5.20)$$

## 6. THE NUMBER OF SPANNING TREES ON $SG_d(n)$ FOR GENERAL $d$

With the above theorems that the numbers of spanning trees on  $SG_2(n)$ ,  $SG_3(n)$  and  $SG_4(n)$  have only simple factors, we observe the following conjecture for the number of spanning trees  $f_d(n) \equiv N_{ST}(SG_d(n))$  on  $SG_d(n)$  for general  $d$ .

**Conjecture 6.1.** *The number of spanning trees on the  $d$ -dimensional Sierpinski gasket  $SG_d(n)$  at stage  $n$  is given by*

$$f_d(n) = 2^{\alpha_d(n)} (d+1)^{\beta_d(n)} (d+3)^{\gamma_d(n)}, \quad (6.1)$$

where the exponents are

$$\alpha_d(n) = \frac{d-1}{2} [(d+1)^n - 1], \quad (6.2)$$

$$\beta_d(n) = \frac{d-1}{2d} [(d+1)^{n+1} + dn + d - 1], \quad (6.3)$$

$$\gamma_d(n) = \frac{d-1}{2d} [(d+1)^n - dn - 1]. \quad (6.4)$$

Notice that these exponents are positive integers when  $d$  is a positive integer and  $n$  is a non-negative integer. It is intriguing that the recursion relations (4.3)–(4.5), (4.15)–(4.17), (5.1)–(5.4) and (5.10)–(5.14) become more and more complicated as  $b$  and  $d$  increase, but the solutions remain simple. We do not know a good reason to explain it by the present method.

$SG_d(0)$  at state  $n = 0$  is a complete graph with  $(d + 1)$  vertices, each of which is adjacent to all of the other vertices. Setting  $n = 0$  into Conjecture 2 is consistent with the expectation that the number of spanning trees on the  $d$ -dimensional Sierpinski gasket  $SG_d(0)$  at stage zero is given by  $f_d(0) = (d + 1)^{d-1}$ .

With the number of vertices for  $SG_d(n)$  given in Eq. (2.6), Conjecture 6.1 leads to following result.

**Conjecture 6.2.** *The asymptotic growth constant for  $SG_d(n)$  is given by*

$$z_{SG_d} = \frac{d - 1}{d(d + 1)} [d \ln 2 + (d + 1) \ln(d + 1) + \ln(d + 3)]. \tag{6.5}$$

Conjectures 6.1 and 6.2 reduce to Theorems 3.1, 5.1, 5.2 and Corollaries 3.2, 5.1, 5.2 respectively when  $d$  is set to two, three and four.

By Eq. (2.4),  $z_{SG_d}$  has the upper bound  $\ln(b_{2d})$ . It is of interest to see how close the value  $z_{SG_d}$  is to this bound. For this purpose, we define the ratio

$$r_{SG_d} = \frac{z_{SG_d}}{\ln(b_{2d})} \tag{6.6}$$

and list the first few values of  $z_{SG_d}$  and  $r_{SG_d}$  in Table I. Our results agree with the observation made in Ref. 14 that  $z_{SG_d}$  increases as the degree  $k = 2d$  increases. Compared with the values  $z_{\mathcal{L}_d}$  for  $d$ -dimensional hypercubic lattice  $\mathcal{L}_d$  which also has  $k = 2d$  in Refs. 14, 34,  $z_{SG_d} < z_{\mathcal{L}_d}$  for all  $d \geq 2$  indicates that  $SG_d$

**Table I. Numerical values of  $z_{SG_d}$  and  $r_{SG_d}$ . The last digits given are rounded off.**

$d$	$D$	$k$	$z_{SG_d}$	$r_{SG_d}$
2	1.585	4	1.048594857	0.862051042
3	2	6	1.569396409	0.928042816
4	2.322	8	1.914853265	0.953722370
5	2.585	10	2.172764568	0.966999152
6	2.807	12	2.378271274	0.974945363
7	3	14	2.548944395	0.980152960
8	3.170	16	2.694814686	0.983785240
9	3.322	18	2.822140640	0.986437211
10	3.459	20	2.935085659	0.988442577

is less densely connected than  $\mathcal{L}_d$ . The ratio  $r_{SG_d}$  increases with dimension and approaches one, but the convergence is not as fast as those for the regular lattices given in Refs. 14, 16.

## ACKNOWLEDGMENTS

The research of S.C.C. was partially supported by the NSC grant NSC-95-2112-M-006-004 and NSC-95-2119-M-002-001. The research of L.C.C was partially supported by the NSC grant NSC-95-2115-M-030-002. L.C.C. thanks Prof. R.-L. Sheu for helpful discussion.

## REFERENCES

1. G. Kirchhoff, Über die Auflösung der Gleichungen, auf welche man bei der Untersuchung der linearen Verteilung galvanischer Ströme geführt wird. *Ann. Phys. Chem.* **72**:497–508 (1847).
2. N. L. Biggs, *Algebraic Graph Theory*, 2nd edn. (Cambridge University Press, Cambridge, 1993).
3. D. J. A. Welsh, *Complexity: Knots, Colourings, and Counting (London Math. Soc. Lecture Notes series 186)*. (Cambridge University Press, Cambridge, 1993).
4. R. Burton and R. Pemantle, Local characteristics, entropy and limit theorems for spanning trees and domino tilings via transfer-impedances. *Ann. Probab.* **21**:1329–1371 (1993).
5. R. Lyons, Asymptotic enumeration of spanning trees, *Combin. Probab. Comput.* **14**:491–522 (2005).
6. H. N. V. Temperley, in: *Combinatorics: Proc. Combinatorial Mathematics*, D. J. A. Welsh and D. R. Woodall, editors. (The Institute of Mathematics and its Applications, Oxford, 1972, pp. 356–357).
7. F.-Y. Wu, Number of spanning trees on a lattice, *J. Phys. A: Math. Gen.* **10**:L113–L115 (1977).
8. C. M. Fortuin and P. W. Kasteleyn, On the random cluster model. I. Introduction and relation to other models. *Physica* **57**:536–564 (1972).
9. F.-Y. Wu, The Potts model, *Rev. Mod. Phys.* **54**:235–268 (1982).
10. D. Dhar, Theoretical studies of self-organized criticality. *Physica A* **369**:29–70 (2006).
11. B. Y. Wu and K.-M. Chao, *Spanning Trees and Optimization Problems* (Chapman & Hall/CRC, Boca Raton, 2004).
12. H. N. V. Temperley, *Combinatorics (London Math. Soc. Lecture Note Series #13)*, T. P. McDonough and V. C. Mavron, editors. (Cambridge University Press, Cambridge, 1974, pp. 202–204).
13. W.-J. Tzeng and F.-Y. Wu, Spanning trees on hypercubic lattices and nonorientable surfaces. *Appl. Math. Lett.* **13**:19–25 (2000).
14. R. Shrock and F.-Y. Wu, Spanning trees on graphs and lattices in  $d$  dimensions. *J. Phys. A: Math. Gen.* **33**:3881–3902 (2000).
15. S.-C. Chang and R. Shrock, Some exact results for spanning trees on lattices. *J. Phys. A: Math. Gen.* **39**:5653–5658 (2006).
16. S.-C. Chang and W. Wang, Spanning trees on lattices and integral identities. *J. Phys. A: Math. Gen.* **39**:10263–10275 (2006).
17. B. B. Mandelbrot, *The Fractal Geometry of Nature* (Freeman, San Francisco, 1982).
18. K. J. Falconer, *Fractal Geometry: Mathematical Foundations and Applications*, 2nd edn. (Wiley, Chichester, 2003).
19. Y. Gefen, B. B. Mandelbrot and A. Aharony, Critical phenomena on fractal lattices. *Phys. Rev. Lett.* **45**:855–858 (1980).

20. Y. Gefen and A. Aharony, Solvable fractal family, and its possible relation to the backbone at percolation. *Phys. Rev. Lett.* **47**:1771–1774 (1981).
21. R. Rammal and G. Toulouse, Spectrum of the Schrödinger equation on a self-similar structure. *Phys. Rev. Lett.* **49**:1194–1197 (1982).
22. S. Alexander, Superconductivity of networks. A percolation approach to the effects of disorder. *Phys. Rev. B* **27**:1541–1557 (1983).
23. E. Domany, S. Alexander, D. Bensimon and L. P. Kadanoff, Solutions to the Schrödinger equation on some fractal lattices. *Phys. Rev. B* **28**:3110–3123 (1983).
24. Y. Gefen, A. Aharony and B. B. Mandelbrot, Phase transitions on fractals: I. Quasi-linear lattices. *J. Phys. A: Math. Gen.* **16**:1267–1278 (1983); Y. Gefen, A. Aharony, Y. Shapir and B. B. Mandelbrot, Phase transitions on fractals: II. Sierpinski gaskets, *ibid.* **17**:435–444 (1984); Y. Gefen, A. Aharony, and B. B. Mandelbrot, Phase transitions on fractals: III. Infinitely ramified lattices, *ibid.* **17**:1277–1289 (1984).
25. R. A. Guyer, Diffusion on the Sierpinski gaskets: A random walker on a fractally structured object. *Phys. Rev. A* **29**:2751–2755 (1984).
26. K. Hattori, T. Hattori and S. Kusuoka, Self-avoiding paths on the pre-Sierpinski gasket. *Probab. Theory Relat. Fields* **84**:1–26 (1990); T. Hattori and S. Kusuoka, The exponent for the mean square displacement of self-avoiding random walk on the Sierpinski gasket, *ibid.* **93**:273–284 (1992).
27. D. Dhar and A. Dhar, Distribution of sizes of erased loops for loop-erased random walks. *Phys. Rev. E* **55**:R2093–R2096 (1997).
28. F. Daerden and C. Vanderzande, Sandpiles on a Sierpinski gasket. *Physica A* **256**:533–546 (1998).
29. D. Dhar, Branched polymers on the Given-Mandelbrot family of fractals. *Phys. Rev. E* **71**:031801 (2005).
30. F. Harary, *Graph Theory* (Addison-Wesley, New York, 1969).
31. B. McKay, Spanning trees in regular graphs. *Eur. J. Combin.* **4**:149–160 (1983).
32. F. Chung and S.-T. Yau, Coverings, heat kernels and spanning trees. *J. Combin.* **6**:163–183 (1999).
33. R. Hilfer and A. Blumen, Renormalisation on Sierpinski-type fractals. *J. Phys. A: Math. Gen.* **17**:L537–L545 (1984).
34. J. L. Felker and R. Lyons, High-precision entropy values for spanning trees in lattices. *J. Phys. A: Math. Gen.* **36**:8361–8365 (2003).